¹ Microfluidic pumps employing surface acoustic waves generated in ZnO ₂ thin films

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ZnO thin film based surface acoustic wave (SAW) devices have been utilized to fabricate 11 microfluidic pumps. The SAW devices were fabricated on nanocrystalline ZnO piezoelectric thin 12 films deposited on Si substrates using rf magnetron sputtering and use a Sezawa wave mode for 13 effective droplet motion. The as-deposited ZnO surface is hydrophilic, with a water contact angle of 14 15 \sim 75°, which prevents droplet pumping. Therefore, the ZnO surface was coated in a self-assembled monolayer of octadecyltrichlorosilane which forms a hydrophobic surface with a water contact 16 angle of $\sim 110^{\circ}$. Liquid droplets between 0.5 and 1 μ l in volume were successfully pumped on the 17 hydrophobic ZnO surface at velocities up to 1 cm s⁻¹. Under acoustic pressure, the water droplet on 18 an hydrophilic surface becomes deformed, and the asymmetry in the contact angle at the trailing and 19 20 leading edges allow the force acting upon the droplet to be calculated. These forces, which increase with input voltage above a threshold level, are found to be in the range of $\sim 100 \ \mu$ N. A pulsed rf 21 signal has also been used to demonstrate precision manipulation of the liquid droplets. Furthermore 22 a SAW device structure is demonstrated in which the ZnO piezoelectric only exists under the input 23 and output transducers. This structure still permits pumping, while avoiding direct contact between 24 the piezoelectric material and the fluid. This is of particular importance for biological 25 26 laboratory-on-a-chip applications. © 2009 American Institute of Physics.

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29 I. INTRODUCTION

Surface acoustic wave (SAW) devices have been in com-30 31 mercial use for more than 60 years, with their main applica-32 tions in communications (filters and oscillators in mobile 33 phone or televisions); automotive sensors (torque and tire 34 pressure), environmental sensors (chemical, vapor, humid-**35** ity), and other industrial and commercial sectors.¹⁻³ Re-36 cently, there has been an increased interest in SAW-based 37 biosensors and microfluidic systems using high performance **38** piezoelectric materials, such as $LiNbO_3$.^{4–6} SAW-based bio-39 chemical sensors normally have high sensitivity and low de-40 tection limits in the order of a few pg/ml. In order to detect 41 the existence of biological species such as cancer cells, pro-42 teins, or DNA, it is essential to handle and manipulate small 43 quantities of fluids (both reagents and specimens) to immo-44 bilize and to bind the target molecules on the surface of the 45 biosensor. Handling of such small volumes of liquid in drop-46 let forms a key challenge.

47 Various microfluidic pumps and mixers have been devel48 oped to control, manipulate, and mix the minute amount of
49 liquid in microliter to picoliter volumes, including devices
50 based on mechanical moving parts (such as oscillating mem51 branes), electric fields applied to liquids, magnetic fields ap-

plied to fluids or inducing phase changes in fluids.^{7,8} SAW- ⁵² based pumps and mixers have distinct advantages, such as ⁵³ simple device structure with no moving parts, low fabrica- ⁵⁴ tion cost, electronic control, high speed, programmable, no ⁵⁵ physical contact between the electrodes and the liquids to be ⁵⁶ manipulated, compactness and high frequency response, and ⁵⁷ the ability to arbitrarily manipulate fluids on a flat surface ⁵⁸ with precision.⁹⁻¹² ⁵⁹

The surface acoustic wave is generated by applying a rf 60 signal to a set of interdigitated transducers (IDTs) which lie 61 on top of a piezoelectric material. The piezoelectric may 62 form the bulk of the substrate, such as lithium niobate 63 $(LiNbO_3)$ or simply be a thin film material on the surface of 64 a substrate, such as a zinc oxide (ZnO) film on a silicon 65 substrate. When the frequency, f, of the rf signal is equal to 66 v_S/p , where v_S is the acoustic velocity of the substrate/ 67 piezoelectric system and p is the periodic spacing of the IDT 68 electrodes, then constructive interference occurs and an in- 69 tense acoustic wave is generated which travels through the 70 piezoelectric substrate. The mode of the acoustic wave is 71 determined by the crystallographic orientation of the piezo- 72 electric material and, in the case of devices using a thin film 73 piezoelectric, the thickness of the piezoelectric layer.¹³ For 74 microfluidic applications, a component of the acoustic wave 75 is required in the direction of propagation, and the so-called 76 Rayleigh mode is commonly employed in which an indi- 77 vidual atom performs elliptical motion in the plane perpen- 78

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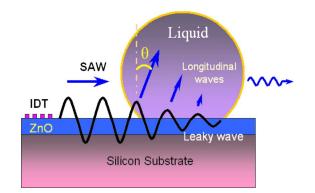


FIG. 1. (Color online) Schematic drawing of interaction between surface acoustic wave and a liquid droplet and defining the Rayleigh angle, θ .

⁷⁹ dicular to the surface and parallel to the direction of propa80 gation. However, the excessive damping of the Rayleigh
81 mode by the liquid means that this mode is considered to be
82 unsuitable for sensing applications.¹⁴⁻¹⁶

 The coupling of the acoustic wave into liquid on the surface of the SAW device, which is required for pumping or mixing, occurs through the excited longitudinal waves propagating into the liquid at an angle called the Rayleigh angle, following the Snell law of diffraction (see Fig. 1).^{17–19} The Rayleigh angle, θ , is defined by

$$\theta = \sin^{-1} \left(\frac{v_L}{v_S} \right),\tag{1}$$

 where v_L is the velocity of the longitudinal wave in the liq- uid. However, the energy and the momentum of the longitu- dinal wave radiated into the liquid are quite useful for liquid pumping and mixing. Indeed, liquid streaming, movement, and even ejection (atomization) have been demonstrated on LiNbO₃ SAW devices upon gradual increase of the wave amplitude.²⁰⁻²³ A net pressure gradient, *P*, forms in the di- rection of the acoustic wave propagation and provides an effective force to drive the liquid, which can be described by²⁴

$$P = \rho_o v_s^2 \left(\frac{\Delta \rho}{\rho_o}\right)^2,\tag{2}$$

1

 in which, ρ_o is the liquid density and $\Delta\rho$ is the slight density change due to the acoustic pressure. Based on the acoustic streaming effect, microfluidic pump and mixer systems,^{25–28} droplet positioning and manipulation systems,^{29,30} atomiza- tion systems,³¹ fluidic dispenser arrays,³⁰ and acoustic ejectors³² have all been proposed and developed.

107 Most of the SAW devices so far have been made from 108 bulk piezoelectric materials, such as LiNbO₃ and quartz. 109 Bulk piezoelectric materials are expensive, and cannot be 110 integrated with control electronics, preventing them from 111 widespread application in microfluidics. On the other hand, 112 thin film ZnO has good piezoelectric properties, a high elec-113 tromechanical coupling coefficient, high sensitivity, and 114 reliability.³³ Furthermore, it can be deposited by such meth-115 ods as rf magnetron sputtering and laser-assisted deposition 116 on a variety of substrates, including silicon, making it the 117 most promising material for integration with electronic con-118 trol circuitry.^{34,35} Such integration is the prerequisite for a 140

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fully automated and digitized microsystem with low cost, ¹¹⁹ fast speed, reduced reagent requirement, and precision con- 120 trol of liquid quantity and position. However, little effort has 121 been made to develop SAW microfluidic systems using pi- 122 ezoelectric thin films such as ZnO or AlN. The authors have 123 recently successfully demonstrated acoustic mixing using the 124 SAW devices made on ZnO thin films, and have obtained a 125 within-droplet streaming velocity of up to 5 cm s⁻¹ using a **126** Sezawa wave (high order Rayleigh mode) in the devices.^{12,36} 127 This paper reports on the effect of a self-assembled mono- 128 layer of octodecyltrichlorosilane (OTS) coating on hydro- 129 phobicity properties of ZnO film surfaces. Consequently, mi- 130 crofluidic pumping is achieved rather than mixing alone. 131 Observation of microcontact angle allows the force acting on 132 water droplets to be determined and general principles for 133 the requirements of hydrophobic surfaces for microfluidic 134 pumps to be established. A SAW structure design is also 135 proposed and demonstrated in which the piezoelectric ZnO 136 thin film is only under the IDTs and not in the pumping 137 region, which is particularly attractive for laboratory-on-a- 138 chip applications. 139

II. EXPERIMENTAL PROCEDURE

The fabrication of the IDT structures has been previ- 141 ously described in detail by the authors in Refs. 12 and 42. 142 Thin films of ZnO up to 6.6 μ m thick were deposited onto 143 Si(100) substrates by rf magnetron sputtering. The IDT trans- 144 ducers were then fabricated from rf magnetron sputtered alu- 145 minum thin films up to 150 nm thick. This was sufficiently 146 thin to avoid significant mass loading effects.³⁷ The resulting 147 IDTs consisted of 30 or 60 pairs of fingers, with a period, p, 148 of 32 μ m, and an aperture (active IDT width) of 4900 μ m. 149 An HP8711A rf network analyzer was used to characterize 150 the SAW devices. The propagation of the SAW was studied 151 through the analysis of the scattering parameters in two con- 152 figurations: S_{11} and S_{22} (reflection) and S_{21} and S_{12} (transmis- 153) sion). The rf signal from an Agilent N9310A signal generator 154 was amplified by a power amplifier up to a peak-to-peak 155 value of 70 V before being fed into the IDTs. The amplitudes 156 of the signals or input voltages were checked using an oscil- 157 loscope (Tektronix TDS). Water droplets with different sizes 158 were obtained using a Micro-Volume Kit micropipette. Drop- 159 let movement was measured using a video camera (Motic 160 MCCamera) using both top and horizontal views. 161

III. RESULTS AND DISCUSSIONS 162

A. Measurement of acoustic force

ZnO is hydrophilic with a contact angle typically at 30° – 164 75°, which is strongly dependent on the surface conditions 165 and light exposure.³⁸ The contact angle increases with the 166 ZnO film thickness as the roughness of the ZnO films in- 167 creases with the thickness. The typical contact angle for a 168 clean surface is between 70° and 80°. The hydrophilic nature 169 of the surface has the effect of preventing microfluidic 170 pumping of water droplets on the SAW device. Instead, 171 streaming only is observed within the droplet, which the au- 172 thors have reported on previously.³⁶ Additionally, the droplet 173 deforms so that the contact angle on the "trailing edge" of 174

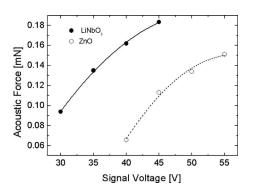


FIG. 2. Acoustic force as a function of rf signal voltage acting on a 10 μ l water droplet for both a LiNbO₃ (Ref. 42) and ZnO SAW device. The LiNbO₃ device uses a Rayleigh mode wave, while the ZnO device uses a Sezawa mode wave. Both devices have 60 finger pairs in the IDT.

175 the drop, θ_t (closest to the driving IDT) is reduced, while that 176 on the leading edge, θ_l , is increased. The acoustic force, F_s , 177 can be calculated from the asymmetry in these contact angles 178 and the droplet size from,^{39–41}

$$F_s = 2R\gamma_{\rm LG}\sin\left(\frac{\theta_t + \theta_l}{2}\right)(\cos \theta_t - \cos \theta_l), \qquad (3)$$

180 where *R* is the radius of the droplet and γ_{LG} is the liquid-gas **181** surface energy.

182 Figure 2 shows the dependence of acoustic force gener-**183** ated on a 10 μ l water drop by the Sezawa wave from the 184 ZnO piezoelectric as a function of rf signal voltage (peak to 185 peak), as well as the acoustic force produced by a Rayleigh 186 wave from a LiNbO₃ SAW IDT that authors have reported.⁴² 187 It was found that, under conditions where a Rayleigh mode 188 wave was induced in the ZnO, then the acoustic force gen-189 erated was insufficient to deform the droplet, whereas when 190 the Sezawa mode wave was generated, significant droplet 191 deformation was observed, indicating a far high acoustic 192 force. This contrasts with the observations for LiNbO₃ SAW **193** devices⁴² where the fundamental Rayleigh mode wave only **194** has sufficient force to deform the droplet. The key difference 195 between the two devices is that, while the ZnO piezoelectric 196 is a thin film material on a silicon substrate, the LiNbO₃ is a 197 bulk piezoelectric wafer. This indicates that different wave 198 modes are required for microfluidic pumps fabricated from 199 bulk and thin film piezoelectrics where, in the latter case, the 200 acoustic wave has a certain proportion confined to the thin 201 film region.

Figure 2 also shows that a threshold voltage, V_{th} , is re-203 quired before droplet deformation is observed. Thereafter, it 204 is found that the force increases with the input voltage, but 205 has a tendency to saturate at high values of voltage. The 206 rolling-off in acoustic force with increasing power is attrib-207 uted to temperature-induced frequency shifts at high 208 power.^{42–44} It is known that the transmitted acoustic wave 209 amplitude, *T*, has the form

210
$$T = \gamma \left(\frac{\sin(\Delta f)}{\Delta f}\right)^2,$$
 (4)

211 where Δf is the deviation in frequency from the resonant **212** frequency and γ is a constant. It may reasonably be assumed **213** that the power fed into the acoustic wave is proportional to

TABLE I. Extracted parameters for the acoustic force extraction.

Material	lpha (m NV ⁻²)	$egin{array}{c} eta\ (\mathrm{V}^{-1}) \end{array}$	$V_{ m th}$ (V)
LiNbO ₃	$(3.28 \pm 0.08) \times 10^{-5}$	0.042 ± 0.006	11±3
ZnO	$(5\pm 2) \times 10^{-4}$	0.055 ± 0.011	27 ± 3

 $(V-V_{\rm th})^2$, and that this has two effects. First, it will increase ²¹⁴ the amplitude of the acoustic wave, and hence the force act-²¹⁵ ing on the droplet. Second, it will cause the system to heat ²¹⁶ up, resulting in a shift in the resonant frequency. Assuming ²¹⁷ that this shift is proportional to the input power also, then ²¹⁸ these effects may be combined in a single expression for the ²¹⁹ acoustic force as a function of input voltage, ²²⁰

$$F_s = \alpha (V - V_{\text{th}})^2 \left(\frac{\sin[\beta(V - V_{\text{th}})]}{\beta(V - V_{\text{th}})}\right)^2,\tag{5}$$

where α is a force-voltage coupling coefficient and β is the 222 constant of proportionality relating the input voltage to the 223 temperature-induced frequency shift. Figure 2 shows the fit- 224 ted curve for each case, and Table I shows the fitting param- 225 eters. It is noticeable that the force-voltage coupling coeffi- 226 cient is greater for ZnO as well as the temperature-induced 227 frequency shift coefficient. The threshold voltage is also 228 much greater in the ZnO device, and is most likely a result of 229 the ZnO being a polycrystalline thin film compared to the 230 bulk single crystal LiNbO₃. The former has a limited piezo- 231 electric effect. The upper limit of the signal voltage for the 232 LiNbO₃ SAW device is 45 V, above which the LiNbO₃ de- 233 vice cracks along the IDTs and fails to operate because of 234 localized heating effect at the electrodes. No such cracking 235 was observed in the ZnO SAW devices up to 70 V, probably 236 due to the higher thermal conductivity of the Si substrate. At 237 extremely high signal voltages (>60 V), a stream of tiny 238 droplets with sizes in the range from a few fl to pl are gen- 239 erated and ejected from the ZnO SAW device. 240

B. Microfluidic pumping on hydrophobic surfaces 241

For droplet manipulation on a channel less flat surface, 242 the surface of the SAW devices was modified to obtain a 243 hydrophobic surface. A self-assembled monolayer (SAM) of 244 octadecyltrichlorosilane (OTS) monolayer was found to be 245 one of the best materials for this purpose. The process to 246 form an OTS SAM layer involved immersing the sample in a 247 toluene solution of OTS (0.01 mM) for 6 h at room tempera- 248 ture and then baking it for 1 h at 125 °C. The IDT electrodes 249 were covered by photoresist (AZ5214E) to avoid Al elec- 250 trode corrosion by the OTS solution. Once the SAM forma- 251 tion was completed, the photoresist was removed by immers- 252 ing the sample in acetone in which the OTS SAM layer is 253 stable. The OTS SAM layer exhibits excellent uniformity 254 and good adhesion to the substrate. The average thickness of 255 the OTS film is ~ 10 nm which induces no measurable 256 acoustic damping. 257

After the ZnO surface was modified by the OTS, the **258** contact angle of the water droplet increased from $\sim 75^{\circ}$ to **259** $\sim 110^{\circ}$, as shown in Fig. 3, largely depending on the process **260**

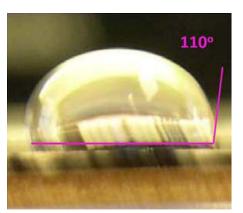


FIG. 3. (Color online) Photos of a droplet on an untreated and an OTS treated ZnO surface. The contact angle increased from \sim 75° to 110° after OTS SAM layer formation.

²⁶¹ conditions. The OTS treated surface of the ZnO SAW is very262 stable and remains hydrophobic with the same contact angle263 even after several weeks.

 Increasing the contact angle of a liquid droplet on a solid surface by making the surface hydrophobic also reduces the work of adhesion, W_s , of the droplet on the surface.⁴⁵ For any liquid on a surface, the work of adhesion is related to the contact angle, θ , and the liquid-gas surface free energy, γ_{LG} , **269** by

$$W_{\rm S} = \gamma_{\rm LG} (1 + \cos \theta), \tag{6}$$

 where γ_{LG} is 72.9 mN m⁻¹ for water at room temperature in air. This is the force per unit length acting at the line of contact between the air, liquid, and solid. Therefore, the total acoustic force that must be applied to the droplet to cause pumping must be greater than the force acting along the length $(2\pi r \sin \theta)$ of this line of contact, F_c , which can be calculated to be

278
$$F_c = \gamma_{\rm LG} (1 + \cos \theta) 2\pi R \sin \theta, \qquad (7)$$

279 where the radius of the droplet, R, may be simply calculated **280** from its volume, V, by

281
$$V = \frac{\pi R^3}{3} (2 + \cos \theta) (1 - \cos \theta)^2.$$
 (8)

282 Figure 4 shows how F_c varies as a function of contact angle **283** for a 10 μ l droplet of water on a surface. It is clear that

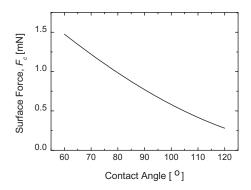


FIG. 4. Variation in minimum acoustic force, F_c , required to move a water droplet of 10 μ l volume on a surface as a function of contact angle calculated from Eqs. (6) and (7).

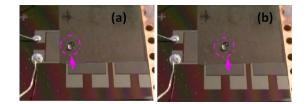


FIG. 5. (Color online) Photos of a 1 μ l droplet movement on an OTS treated ZnO SAW device before (a) and after (b) being driven by a rf signal at a frequency of 178.7 MHz.

reducing the contact angle from $\sim 75^{\circ}$ to 110° results in a ²⁸⁴ reduction in the work of adhesion by almost a magnitude, 285 from well in excess of 1.5 mN to the order of $\sim 400 \ \mu N$. 286 This is the force that would be required to detach the droplet 287 from the surface completely. The force required to simply 288 move the droplet will be some fraction of the work of adhe- 289 sion. It has already been shown in Sec. III A that the acoustic 290 force that can be generated using the SAW devices in this 291 study is ~100 μ N for a 10 μ l droplet. The exact mecha- 292 nism of droplet motion is not clear yet, but two possible 293 mechanisms are suggested here. First, under such a strong 294 acoustic force the droplet deforms with an increased leading 295 edge contact angle and decreased trailing edge contact angle, 296 as shown in a previous publication by the authors.⁴² The 297 increased surface energy of the droplet on the hydrophobic 298 surface makes the contact angle of the trailing edge return to 299 its original value, which reduces the contact area, while the 300 conservation of the work of adhesion of the droplet makes 301 the leading angle decrease by expanding the droplet area a 302 little in front, hence moving the droplet forward a little. By 303 applying a "continuous" rf wave, a constant droplet velocity 304 appears to result. This pumping process is similar to that of 305 electrowetting,^{46,47} but in a reversed order (i.e., the droplet 306 moves in the increased contact angle direction under the 307 acoustic force). The second possible mechanism is that, un- 308 der the application of the rf wave, the deformed droplet vi- 309 brates, and may expand a little in forward direction owing to 310 the acoustic force, hence moving the droplet forward. 311

Figure 5 shows photos of a 1 μ l droplet before and after 312 applying a rf signal at 178.7 MHz to the ZnO SAW device 313 with a peak-to-peak voltage of 40 V for a duration of 2 s. 314 The water droplet is observed to have moved a distance of 315 ~4 mm. Figure 6 shows the dependence of droplet velocity 316 as a function of rf signal voltage. Unsurprisingly, this graph 317

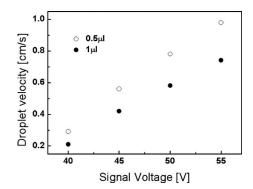


FIG. 6. Sezawa wave driven droplet velocity as a function of rf signal voltage applied on the SAW for a 5.5 μ m ZnO SAW device.

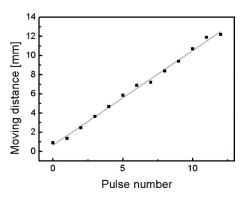


FIG. 7. Total distance moved by a 0.5 μ l water droplet as a function of pulse number for a rf signal of 40 V and a pulse width of 100 ms.

³¹⁸ shows a similar form to that found for acoustic force as a 319 function of input voltage (Fig. 2). Smaller droplets are found 320 to have a higher velocity for the same input voltage due to 321 their reduced work of adhesion. The droplet velocities **322** achieved on ZnO are about half those reported previously by **323** the authors on LiNbO₃,⁴² but are larger than those of most **324** micropumps and are sufficient for microfluidic applications. Practically, a pulsed rf signal is normally used to control 325 326 the droplet motion, as this affords more precise control of the 327 position and moving distance while also suppressing acoustic 328 heating. Figure 7 shows the dependence of the moving dis-**329** tance of a 0.5 μ l water droplet as a function of pulse number 330 applied. The rf signal voltage is 40 V and the pulse width is 331 100 ms. The total distance moved increases linearly with the 332 number of pulses applied. The average distance moved per **333** pulse is nearly 100 μ m, although this is strongly dependent 334 on the signal voltage and the droplet size.

335 C. Island structure SAW devices

Although ZnO is not biologically toxic, it is very reac-337 tive with acids or bases⁴⁸ and will dissolve or recrystallize if 338 exposed to water or a humid environment.⁴⁹ Therefore, ZnO 339 SAW devices with no surface protection are not suitable for 340 laboratory-on-a-chip applications. Surface coating is consid-341 ered to be one way of protecting ZnO thin films from deg-342 radation. However, in this work a ZnO SAW structure is 343 proposed which can avoid a direct contact between the ZnO 344 active layer and the fluid being pumped altogether.

Since the acoustic wave in a thin film ZnO SAW device 346 is strongly affected by wave propagation in the Si substrate, 347 it is proposed that the two functions of SAW generation and 348 propagation can be separated. The proposed device structure 349 is shown in Fig. 8. The ZnO is patterned into islands so that 350 it only exists underneath the IDTs where piezoelectric mate-351 rial is required for SAW generation or sensing. The base 352 silicon substrate is then used for SAW propagation.

Figure 9 is a comparison of both Rayleigh and Sezawa Transmission spectra of SAW devices with and without The SAW devices were made on the The SAW device on the the SAW device on the transmission The SAW device on a 1.2 μ m ZnO film, the abso-

Standard SAW device

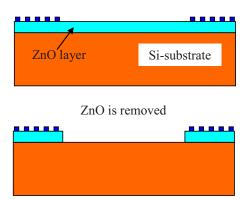


FIG. 8. (Color online) Schematic drawing of the proposed SAW structure with ZnO on the wave path being removed compared with the standard SAW device structure where the ZnO is not patterned.

lute value of the transmitted signal of the Rayleigh and ³⁶¹ Sezawa mode waves decreases after removing the ZnO in the 362 propagation region. However, a signal is still clearly ob- 363 served in both cases that is well above the level of system 364 noise. 365

The resonant SAW frequency is determined by the IDT 366 periodicity and the acoustic speeds in the substrate and thin 367 film overlayer in the region of the IDTs-it is not strongly 368 affected by the active layer in the propagation region. For a 369 layered structure, the longitudinal acoustic waves travel with 370 a certain proportion of the waves existing in each of the thin 371 film surface layer and the bulk substrate. Since the wave- 372 length in this case is $\sim 32 \ \mu m$, which is much greater than 373 the thickness of the 1.2 μ m ZnO thin film, it might be ex- 374 pected that the acoustic energy mainly exists within the sili- 375 con substrate, and that the removal of such a thin layer in the 376 propagation region would not affect the transmission signal 377 significantly. However, the lower acoustic velocity of the 378 ZnO ($\sim 2700 \text{ m s}^{-1}$) relative to the silicon substrate 379 $(\sim 4680 \text{ m s}^{-1})$ means that there will be a tendency for en- **380** hanced wave trapping.^{42,50} This is particularly true for the **381** Sezawa wave mode, which is known to be more confined to 382 the surface. Further increase in the ZnO thickness means a 383 higher proportion of acoustic waves, and hence energy, is 384 trapped within the active layer. This explains the significant 385 decrease in transmission signal observed upon removal of 386 the ZnO layer—an effect that is particularly pronounced for 387 the Sezawa mode wave. 388

This reduction in transmitted signal also results in a re- 389 duced acoustic force that can act upon microfluidic droplets. 390 Figure 10 shows the streaming velocity that occurs within a 391 droplet of water that is exposed to the acoustic wave, but 392 where the surface is hydrophilic so that no pumping motion 393 occurs (the methodology for conducting this experiment has 394 been previously described by the authors in Ref. 42). Stream-395 ing velocities are obtained that are $\sim 60\%$ lower compared to 396 devices with a continuous ZnO layer, but are still significant, 397 and indicates that these devices can still be effectively used 398 for microfluidic applications. 399

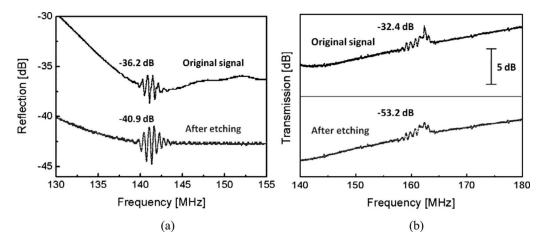


FIG. 9. Comparison of transmission spectra for SAW devices before and after removal of the ZnO from the propagation area using (a) the Rayleigh wave from a 1.2 μ m thick ZnO thin film and (b) the Sezawa wave from a 6.6 μ m thick ZnO thin film.

⁴⁰⁰ D. Acoustic heating

401 Acoustic heating is a common phenomenon in SAW-402 based devices, and will affect the performance of acoustic 403 streaming and pumping significantly.^{51,52} First, increasing 404 temperature shifts the resonant frequency of a SAW device, 405 which is normally taken as a constant in SAW microfluidics. 406 Second, it introduces errors in sensing using the SAW de-407 vices, as changes in transmission amplitude are no longer 408 solely a function of mass loading. Third, the heat may dam-409 age the integrity of biological substances to be tested. In 410 order to quantify this effect, the surface temperature of the 411 SAW devices have been investigated by using a thermo-412 couple to measure the temperature change on the ZnO film 413 surface. Measurements were performed without any liquid 414 on the SAW device surface.

Figure 11 shows the temperature change in the ZnO film Figure 11 shows the temperature change in the ZnO film Figure (5 mm away from the IDT) measured as a function of the apparent and peak-to-peak voltage of the input rf signal. It is reases with increase in the voltage and duration of the rf figure and the maximum temperature recorded was ~ 140 °C for the highest signal voltage of 60 V. Hence, the previous solution of bubble formation and evaporation of the wa-

> 2.0 [Sup 1.5 1.5 1.5 0.5 0.0 1.2µm "R" N=60 0.5 0.0 1.2µm "R" N=60 0.5 1.2µm "R" N=60 0.5 5 Signal voltage [V]

FIG. 10. (Color online) Streaming velocity induced by acoustic waves from the SAW devices where there is no ZnO in the propagation region. Results for both Sezawa (s) and Rayleigh (r) mode devices are presented. ter droplets during ZnO SAW microfluidic studies can be 423 attributed to the acoustic heating effect, although the pres- 424 ence of liquid on the SAW surface will undoubtedly result in 425 a lower temperature during microfluidic operation than those 426 measured here. 427

The acoustic heating can be significantly reduced by us- 428 ing a pulsed rf signal. The surface temperature of the device 429 was found to remain constant at room temperature with a 430 pulse width up to 200 ms and peak-to-peak magnitude of 45 431 V owing to heat dissipation during the off period of the 432 pulsed rf signal, and a slight increase in temperature up to 433 \sim 40 °C was observed when the pulse width increased to 434 600 ms. A pulsed rf signal also has the advantage of inducing 435 effective mixing because of the nonregular streaming pat- 436 terns, and precise control of droplet motion as discussed in 437 Sec. III B. 438

IV. CONCLUSIONS

ZnO thin film based SAW micropumps have been fabri- 440 cated based on low cost, high quality, *c*-axis oriented ZnO 441 thin films deposited on silicon substrates using a rf magne- 442 tron sputtering system. The results showed that droplet ma- 443 nipulation on the ZnO device depends on the SAW ampli- 444

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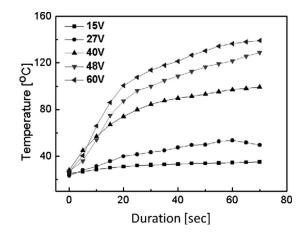


FIG. 11. Temperature change as a function of time with rf signal voltage as a parameter. The temperature rises rapidly at the initial 20 s, and then increases slowly thereafter.

⁴⁴⁵ tude, wave mode, droplet size, and surface hydrophobicity of 446 the substrate. Acoustic forces $\sim 100 \ \mu N$ were produced, and 447 it has been shown that this is sufficient to induce pumping of **448** 10 μ l water droplets where the surface contact angle is sig-449 nificantly greater than 100°. This can be achieved by using 450 OTS to modify the ZnO surface to form a hydrophobic coat-**451** ing with a water contact angle of $\sim 110^{\circ}$. Pumping velocities **452** approaching 1 cm s^{-1} were measured. A pulsed rf signal has 453 also been used to demonstrate precision manipulation of the 454 liquid droplets. Furthermore a SAW device on ZnO island 455 structure has been proposed, fabricated and microfluidic op-456 eration verified. This structure avoids a direct contact be-457 tween the piezoelectric material and microfluidic substances, 458 and hence is promising for laboratory.-on-a-chip applica-**459** tions.

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- #1 Au: Please check changes in Ref. 5.
- #2 Au: Please update Ref. 42.